

Pseudo-differential operators

Summer School - From kinetic equations to statistical mechanics

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Notation

Japanese bracket: $\langle \zeta \rangle := (1 + |\zeta|^2)^{\frac{1}{2}}$

C_b^∞ : space of smooth functions bounded together with all their derivatives

Fourier transform: $\hat{u}(\xi) = \int e^{-ix \cdot \xi} u(x) dx$

Multiindices: for $x \in \mathbb{R}^n$, $\alpha \in \mathbb{N}^n$, $x^\alpha := x_1^{\alpha_1} \cdots x_n^{\alpha_n}$,
 $|\alpha| = \alpha_1 + \cdots + \alpha_n$ and $\partial_x^\alpha = \partial_{x_1}^{\alpha_1} \cdots \partial_{x_n}^{\alpha_n}$

$D_x = i^{-1} \nabla_x$ and $D_x^\alpha = i^{-|\alpha|} \partial_x^\alpha$

Schwartz space: $\mathcal{S} \ni u$ if $x^\alpha \partial_x^\beta u \in L^\infty$ for all α, β

Temperate distributions: \mathcal{S}'

$\langle T, u \rangle =$ bilinear pairing between $T \in \mathcal{S}'$ and $u \in \mathcal{S}$, “ $\int Tu$ ”

$(T, u) = \overline{\langle T, \bar{u} \rangle} =$ sesquilinear pairing between T and u , “ $\int \bar{T}u$ ”

Ψ DO (part 1): Symbol classes

Definition. Let $m, \mu \in \mathbb{R}$: $a \in S^{\mu, m}$ iff for all $\alpha, \beta \in \mathbb{N}^n$

$$|\partial_x^\alpha \partial_\xi^\beta a(x, \xi)| \leq C_{\alpha\beta} \langle x \rangle^\mu \langle \xi \rangle^{m-|\beta|}, \quad (x, \xi) \in \mathbb{R}^{2n}.$$

Remark. When $\mu = 0$, $S^{0, m} =: S^m$ is the most standard class.

Example. If $m \in \mathbb{N}$ and $a_\alpha \in C_b^\infty(\mathbb{R}^n)$

$$a(x, \xi) := \sum_{|\alpha| \leq m} a_\alpha(x) \xi^\alpha \in S^m.$$

Facts.

- ▶ If $m \leq m'$ and $\mu \leq \mu'$, $S^{\mu, m}$ is contained in $S^{\mu', m'}$.
- ▶ If $a \in S^{\mu_1, m_1}$ and $b \in S^{\mu_2, m_2}$ then $ab \in S^{\mu_1 + \mu_2, m_1 + m_2}$.

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Remark. A linear (or conjugate linear) map B from S^{μ_1,m_1} to S^{μ_2,m_2} is continuous if: $\forall N_2 \in \mathbb{N} \exists C > 0$ and $\exists N_1 \in \mathbb{N}$ s.t.

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Definition. If $a \in S^{\mu,m}$, $u \in \mathcal{S}$ and $h \in (0, 1]$,

$$a(x, hD)u(x) = Op_h(a)u(x) := (2\pi)^{-n} \int e^{ix \cdot \xi} a(x, h\xi) \hat{u}(\xi) d\xi.$$

Remark. The case $h = 1$ is relevant.

Where does this definition come from ?

More generally

$$a(x)D^\alpha u(x) = (2\pi)^{-n} \int e^{ix \cdot \xi} a(x)\xi^\alpha \hat{u}(\xi) d\xi$$

so if $p(x, \xi) := \sum_{|\alpha| \leq m} a_\alpha(x)\xi^\alpha$,

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Pseudo-differential operators extend this formula to a larger class of symbols p .

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Proposition. For each h , the map $(a, u) \mapsto Op_h(a)u$ is continuous between $S^{\mu, m} \times \mathcal{S}$ and \mathcal{S} .

Proof.

$$|Op_h(a)u(x)| \lesssim_h \int \underbrace{\langle x \rangle^\mu \langle \xi \rangle^m}_{|a| \lesssim_h} \underbrace{\langle \xi \rangle^{-m-n-1}}_{|\hat{u}| \lesssim} d\xi \lesssim_h \langle x \rangle^\mu$$

and then more generally, for all $\alpha, \beta \in \mathbb{N}^n$

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- **Composition.** For $a \in S^{\mu,m}$, $b \in S^{\mu',m'}$

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where $c(h) \sim \sum h^j c_j$ in $S^{\mu+\mu',m+m'}$, i.e. for all N

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- **Adjoint.** For $a \in S^{\mu,m}$,

$$(Op_h(a)v, u) = (v, Op_h(a^*(h))u), \quad u, v \in \mathcal{S}$$

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for all $h \in (0, 1]$, all $a \in S^{0,0}$ and all $u \in \mathcal{S}$ (or $u \in L^2$).

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Interlude: Some references

Non-semiclassical

S. Alinhac-P. Gérard, *Op. pseudo. diff. et th. de Nash Moser*

G.B. Folland, *Harmonic analysis in phase space*

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Semiclassical

M. Dimassi-J. Sjöstrand, *Spectral asymptotics in the semi-class. limit*

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D. Robert, *Autour de l'approximation semi-classique*

M. Zworski, *Semiclassical analysis*

Application: Hörmander's Theorem on hypoelliptic operators

Consider a differential operator on \mathbb{R}^n of the form

$$P = -\sum_{j=1}^r L_j^2 + L_0 + V,$$

with L_0, L_1, \dots, L_r real vector fields and V a smooth function.

Assumptions There exists N such that for all x ,

$$\{L_k, [L_{k_1}, L_{k_2}], \dots, [L_{k_1}, [L_{k_2}, \dots [L_{k_{N-1}}, L_{k_N}]]\}$$

span $T_x \mathbb{R}^n$. Assume also V and the L_k 's coefficients are C_b^∞ .

Example (local):

$$P = -\frac{\partial^2}{\partial x_1^2} - \sum_{j=2}^n x_1^{m_j} \frac{\partial}{\partial x_j}, \quad 0 \leq m_2 < \dots < m_n.$$

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Theorem. *There exists $\epsilon \in (0, 1]$ ($\epsilon = 2^{-N}$) such that for each $\chi \in C_0^\infty(\mathbb{R}^n)$*

$$\|\chi u\|_{(\epsilon)} \lesssim_\chi \|Pu\|_{L^2} + \|u\|_{L^2}$$

for all $u \in \mathcal{S}(\mathbb{R}^n)$.

Meaning. Even if P is not elliptic, having Pu (and u) in L^2 still gives a control on some (fractional) derivatives of u .

Interests.

- ▶ Can be used to prove the hypoellipticity of P :
 $Pu \in C^\infty \Rightarrow u \in C^\infty$
- ▶ On a compact manifold, it can be used to show that the resolvent of P is compact \Rightarrow discrete spectrum.

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Proof: First step

Rewrite

$$P = \sum_{j=1}^r L_j^* L_j + \tilde{L}_0 + V$$

with $\tilde{L}_0 = L_0 + \sum c_j L_j$ if $L_j^* = -L_j - c_j$. Note that

$$\tilde{L}_0^* = -\tilde{L}_0 + c_0, \quad c_0 \in C_b^\infty(\mathbb{R}^n).$$

Now

$$\operatorname{Re}(Pu, u) = \sum_{j=1}^r \|L_j u\|^2 + \operatorname{Re}(Vu, u) + \frac{((\tilde{L}_0 + \tilde{L}_0^*)u, u)}{2}$$

Thus

$$\sum_{j=1}^r \|L_j u\|^2 \leq (Pu, u) + C\|u\|^2 \lesssim \|Pu\|^2 + \|u\|^2.$$

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Let us decompose $P = P' + iP''$ with

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Up to adding $C \gg 1$ to P we may assume that $P' \geq 0$ (see later). It also suffices to show

$$\|P'' u\|_{(-1/2)} \lesssim \|Pu\| + \|u\|.$$

$$\begin{aligned}
\|P''u\|_{(-\frac{1}{2})}^2 &= \|\langle D \rangle^{-1/2} P''u\|^2 = \left(\underbrace{P''}_{=(P-P')/i} u, \underbrace{\langle D \rangle^{-1} P'' u}_{=: A \in \text{Op}(S^0)} \right) \\
&\leq \|Pu\| \underbrace{\|Au\|}_{\lesssim \|u\|} + |(P'u, Au)|
\end{aligned}$$

Then, by Cauchy-Schwarz for the sesquilinear form $(P'u, w)$,

$$|(P'u, Au)| \leq \underbrace{(P'u, u)^{\frac{1}{2}}}_{(\text{Re}(Pu, u))^{\frac{1}{2}}} (P'Au, Au)^{\frac{1}{2}}$$

while $2(P'Au, Au)$ reads $(PAu, Au) + (Au, PAu)$ i.e.

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we get

$$\begin{aligned} ([P, A]u, Au) &= \sum_1^r ([L_j^*, A]L_j u, Au) + ([L_j, A]u, L_j Au) + O(\|u\|^2) \\ &= \sum_1^r ([L_j^*, A]L_j u, Au) + ([L_j, A]u, AL_j u) \\ &\quad + \underbrace{\|[L_j, A]u\|^2}_{\lesssim \|u\|^2} + O(\|u\|^2) \\ &\lesssim \|Pu\|^2 + \|u\|^2. \end{aligned}$$

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This ends Step 2.

Proof: the general induction step

We let $s_k = 2^{1-k} - 1$ for $k \geq 1$ ($s_1 = 0$, $s_2 = -\frac{1}{2}$). We have seen

$$\|L_j u\|_{(s_1)} \lesssim \|Pu\| + \|u\| \quad (1 \leq j \leq r)$$

$$\|\tilde{L}_0 u\|_{(s_2)} \lesssim \|Pu\| + \|u\|$$

One defines next

$$\mathcal{Q}_1 = \{iL_j \mid j = 1, \dots, r\} \quad \mathcal{Q}_2 = \underbrace{\{i\tilde{L}_0\}}_{\text{or eq. } \{P''\}} \cup \{i[Q, Q'] \mid Q, Q' \in \mathcal{Q}_1\}$$

and, for $k \geq 2$,

$$\mathcal{Q}_{k+1} = \{i[Q, Q'] \mid (Q, Q') \in (\mathcal{Q}_k \times \mathcal{Q}_1) \cup (\mathcal{Q}_{k-1} \times \mathcal{Q}_2)\}$$

Remark. For each $Q_k \in \mathcal{Q}_k$, $Q_k^* - Q_k \in Op(S^0)$.

Induction assumption: $\|Q_k u\|_{(s_k)} \lesssim \|Pu\| + \|u\|$

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Induction assumption: $\|Q_k u\|_{(s_k)} \lesssim \|Pu\| + \|u\|$

Proof: the general induction step

We let $s_k = 2^{1-k} - 1$ for $k \geq 1$ ($s_1 = 0$, $s_2 = -\frac{1}{2}$). We have seen

$$\|L_j u\|_{(s_1)} \lesssim \|Pu\| + \|u\| \quad (1 \leq j \leq r)$$

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One defines next

$$\mathcal{Q}_1 = \{iL_j \mid j = 1, \dots, r\} \quad \mathcal{Q}_2 = \underbrace{\{i\tilde{L}_0\}}_{\text{or eq. } \{P''\}} \cup \{i[Q, Q'] \mid Q, Q' \in \mathcal{Q}_1\}$$

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 &\quad - (Q_1 u, A Q_k u) - (Q_1 u, C u) \\
 &= O(\|Q_1 u\|_{(S_1)} \|Q_k u\|_{(S_k)} + \|Q_1 u\|_{(S_1)} \|u\| \\
 &\quad + \|Q_k u\|_{(S_k)} \|u\|)
 \end{aligned}$$

Proof: estimate on $i[Q, Q']$ with $Q \in \mathcal{Q}_k$ and $Q' \in \mathcal{Q}_1$

$$\begin{aligned}
 \|[Q_1, Q_k]u\|_{(S_{k+1})}^2 &= (\langle D \rangle^{S_{k+1}} [Q_1, Q_k]u, \langle D \rangle^{S_{k+1}} [Q_1, Q_k]u) \\
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 &= (Q_k u, Q_1^* A u) - (Q_1 u, Q_k^* A u) \\
 &= (Q_k u, A Q_1 u) + (Q_k u, \underbrace{((Q_1^* - Q_1)A + [Q_1, A]) u}_{B \in \text{Op}(S^{S_k})}) \\
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Proof: estimate on $i[Q, Q']$ with $Q \in \mathcal{Q}_{k-2}$ and $Q' = P''$

$$\begin{aligned} \|[P'', Q_{k-1}]u\|_{(s_{k+1})}^2 &= ([P'', Q_{k-1}]u, Au) \quad (\text{with } A \in \text{Op}(S^{s_k})) \\ &= i(Q_{k-1}u, (P' - P)Au) + i((P - P')u, Q_{k-1}^*Au) \end{aligned}$$

where

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$$\begin{aligned} |(Q_{k-1}u, PAu)| &\lesssim \|Q_{k-1}u\|_{(s_{k-1})} \|PAu\|_{(-s_{k-1})} \\ &\lesssim \|Q_{k-1}u\|_{(s_{k-1})} (\|APu\|_{(-s_{k-1})} + \|[P, A]u\|_{(-s_{k-1})}) \\ &\lesssim \|Q_{k-1}u\|_{(s_{k-1})} (\|Pu\|_{(s_k - s_{k-1} \leq 0)} + \|[P, A]u\|_{(-s_{k-1})}) \end{aligned}$$

in which $[P, A] = -2 \sum_1^r \text{Op}(S^{s_k})L_j + \text{Op}(S^{s_k})$ so

$$\|[P, A]u\|_{(-s_{k-1})} \lesssim \sum_1^r \|L_j u\|_{(s_k - s_{k-1} \leq 0)} + \|u\|_{(s_k - s_{k-1} \leq 0)}$$

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Proof: estimate on $i[Q, Q']$ with $Q \in \mathcal{Q}_{k-2}$ and $Q' = P''$

$$\begin{aligned} \|[P'', Q_{k-1}]u\|_{(s_{k+1})}^2 &= ([P'', Q_{k-1}]u, Au) \quad (\text{with } A \in \text{Op}(S^{s_k})) \\ &= i(Q_{k-1}u, (P' - P)Au) + i((P - P')u, Q_{k-1}^*Au) \end{aligned}$$

where

$$\begin{aligned} (Pu, Q_{k-1}^*Au) &= (Pu, AQ_{k-1}u) + O(\|Pu\| \|u\|) \\ &= O(\|Pu\| \|Q_{k-1}u\|_{(s_k)}) + O(\|Pu\| \|u\|) \\ &= O(\|Pu\| \|Q_{k-1}u\|_{(s_{k-1})}) + O(\|Pu\| \|u\|) \end{aligned}$$

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Conclusion of the proof

Denote by $(V_\ell)_{1 \leq \ell \leq M}$ the (finite) collection of vector fields

$$\{L_k, [L_{k_1}, L_{k_2}], \dots, [L_{k_1}, [L_{k_2}, \dots [L_{k_{N-1}}, L_{k_N}]]\}.$$

Near any given x_0 , one can find n of them $V_{\ell_1}, \dots, V_{\ell_n}$ of rank n (as vectors of $T_x \mathbb{R}^n$, x close to x_0) and a matrix of smooth functions $(a_{ij}(x))_{1 \leq i, j \leq n}$ such that

$$\frac{\partial}{\partial x_i} = \sum_{j=1}^n a_{ij}(x) V_{\ell_j}, \quad |x - x_0| \ll 1.$$

If $\chi \in C_0^\infty$ is supported near x_0 , we infer

$$\|\chi \partial_{x_i} u\|_{(\epsilon-1)} \lesssim \sum_{j=1}^n \|V_{\ell_j} u\|_{(\epsilon-1)} \lesssim \|Pu\| + \|u\|$$

hence

$$\|\chi u\|_{(\epsilon)} \lesssim \sum \|\partial_{x_i}(\chi u)\|_{(\epsilon-1)} + \|u\|_{(0)} \lesssim \underbrace{\|(\nabla \chi)u\|_{(\epsilon-1)}}_{\lesssim \|u\|_{(0)} = \|u\|} + \|Pu\| + \|u\| \quad \square$$

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$$\|\chi \partial_{x_i} u\|_{(\epsilon-1)} \lesssim \sum_{j=1}^n \|V_{\ell_j} u\|_{(\epsilon-1)} \lesssim \|Pu\| + \|u\|$$

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Conclusion of the proof

Denote by $(V_\ell)_{1 \leq \ell \leq M}$ the (finite) collection of vector fields

$$\{L_k, [L_{k_1}, L_{k_2}], \dots, [L_{k_1}, [L_{k_2}, \dots [L_{k_{N-1}}, L_{k_N}]]\}.$$

Near any given x_0 , one can find n of them $V_{\ell_1}, \dots, V_{\ell_n}$ of rank n (as vectors of $T_x \mathbb{R}^n$, x close to x_0) and a matrix of smooth functions $(a_{ij}(x))_{1 \leq i, j \leq n}$ such that

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We let $h = 1$. Pick $a_\epsilon \in C_0^\infty$ such that $a_\epsilon \rightarrow a$ in $S^{0,0}$. Observe that

$$\begin{aligned}Op_h(a_\epsilon)u(x) &= (2\pi)^{-n} \int \int e^{i(x-y)\cdot\xi} a_\epsilon(x, \xi) u(y) dy d\xi \\ &= \int K_\epsilon(x, y) u(y) dy\end{aligned}$$

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By density, we may work with symbols in C_0^∞ . Let A_h (resp. B_h)

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Ψ DO (part 1): Proof of the composition formula

The remainder terms associated to the previous expression are of the form, with $|\alpha| = N$

$$(2\pi)^{-n} \int_0^1 (1-t)^{N-1} \int e^{i\eta \cdot z} \hat{a}(x, z) (\partial_x^\alpha b)(x + thz, \eta) \frac{(hz)^\alpha}{\alpha!} dz dt$$

Dropping the integration in t , we get for each $t > 0$

$$h^N \int e^{i\eta \cdot z} \underbrace{z^\alpha \hat{a}(x, z)}_{\widehat{D_\xi^\alpha a}(x, z)} (\partial_x^\alpha b)(x + thz, \eta) dz =$$
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Ψ DO (part 1): Proof of the composition formula

To complete the proof, we wish to gain additional **polynomial** decay in z and ξ by integrating by part using

$$\frac{1}{(|z|^2 + 1)^M} (1 - \Delta_\xi)^M e^{i(\eta - \xi) \cdot z} = e^{i(\eta - \xi) \cdot z}$$

and

$$\frac{1}{(|\eta - \xi|^2 + 1)^M} (1 - \Delta_z)^M e^{i(\eta - \xi) \cdot z} = e^{i(\eta - \xi) \cdot z}$$

so we end up with integrals of the form (if we view a and b as symbols in $S^{\mu, m}$ and $S^{\mu', m'}$, as will eventually be the case)

$$\iint O(\langle z \rangle^{-2M}) O(\langle \xi - \eta \rangle^{-2M}) O(\langle x \rangle^\mu \langle \xi \rangle^{m-N}) O(\langle x + thz \rangle^{\mu'} \langle \eta \rangle^{m'}) d\xi dz$$

i.e.

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Ψ DO (part 1): Proof of the composition formula

To complete the proof, we wish to gain additional **polynomial** decay in z and ξ by integrating by part using

$$\frac{1}{(|z|^2 + 1)^M} (1 - \Delta_\xi)^M e^{i(\eta - \xi) \cdot z} = e^{i(\eta - \xi) \cdot z}$$

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so we end up with integrals of the form (if we view a and b as symbols in $S^{\mu, m}$ and $S^{\mu', m'}$, as will eventually be the case)

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Ψ DO (part 2): More general classes

Definition. A positive function m on \mathbb{R}^{2n} is an **order function** if there exist C and M such that

$$m(x, \xi) \leq Cm(y, \eta)(1 + |x - y| + |\xi - \eta|)^M.$$

Examples. $m(x, \xi) = \langle x \rangle^\mu \langle \xi \rangle^\nu$
In $\mathbb{R}^3 \times \mathbb{R}^3$

$$m(v, \eta) = \langle v \rangle^\gamma (1 + |\eta|^2 + |\eta \wedge v|^2 + |v^2|)^s$$

with $0 < s < 1$ and $\gamma > -3$ (related to the linearized Boltzmann operator)

Definition. Let $g_0 = dx^2 + d\xi^2$ the standard Euclidean metric on \mathbb{R}^{2n} . If m is an order function, $\mathcal{S}(m, g_0)$ is the space of smooth functions a such that

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There is a general notion of Riemannian metrics for pseudo-differential calculus (Hörmander). We just give a few examples

$$g_{\text{class}} = dx^2 + \frac{d\xi^2}{\langle \xi \rangle^2}, \quad g_{\text{scat}} = \frac{dx^2}{\langle x \rangle^2} + \frac{d\xi^2}{\langle \xi \rangle^2}, \quad g_{\rho, \delta} = \frac{dx^2}{\langle \xi \rangle^{-2\delta}} + \frac{d\xi^2}{\langle \xi \rangle^{2\rho}}$$

($0 \leq \delta \leq \rho \leq 1$ and $\delta < 1$) and some related symbol classes

$$a \in S(\langle \xi \rangle^m, g_{\text{class}}) = S^m \Leftrightarrow \partial_x^\alpha \partial_\xi^\beta a(x, \xi) = O(\langle \xi \rangle^{m-|\beta|})$$

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$$g_{\text{class}} = dx^2 + \frac{d\xi^2}{\langle \xi \rangle^2}, \quad g_{\text{scat}} = \frac{dx^2}{\langle x \rangle^2} + \frac{d\xi^2}{\langle \xi \rangle^2}, \quad g_{\rho, \delta} = \frac{dx^2}{\langle \xi \rangle^{-2\delta}} + \frac{d\xi^2}{\langle \xi \rangle^{2\rho}}$$

($0 \leq \delta \leq \rho \leq 1$ and $\delta < 1$) and some related symbol classes

$$a \in S(\langle \xi \rangle^m, g_{\text{class}}) = S^m \Leftrightarrow \partial_x^\alpha \partial_\xi^\beta a(x, \xi) = O(\langle \xi \rangle^{m-|\beta|})$$

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Ψ DO (part 2): Other quantizations

The **Weyl quantization** of a symbol a is defined by

$$Op_h^W(a)u(x) = (2\pi)^{-n} \int \int e^{i(x-y)\cdot\xi} a\left(\frac{x+y}{2}, h\xi\right) u(y) dy d\xi.$$

More generally, one sometimes considers for $t \in [0, 1]$

$$Op_{h,t}(a)u(x) = (2\pi)^{-n} \int \int e^{i(x-y)\cdot\xi} a(tx + (1-t)y, h\xi) u(y) dy d\xi.$$

All these quantizations can be written using Op_h , i.e.

$$Op_{h,t}(a) = Op_h(a_t(h)),$$

for some $(a_{t,h})_{h \in (0,1]}$ in the same class as a .

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Ψ DO (part 2): Symbolic calculus and L^2 bounds revisited

Theorem. Let m_1, m_2 be order functions and $a \in S(m_1, g_0)$, $b \in S(m_2, g_0)$.

- ▶ **Composition.** $Op_h^W(a)Op_h^W(b) = Op_h^W(c(h))$ with $c(h) \in S(m_1 m_2, g_0)$ of the form

$$c(h)(x, \xi) \sim \sum_j \left(\frac{h}{2i}\right)^j (\partial_y \cdot \partial_\xi - \partial_x \cdot \partial_\eta)^j a(x, \xi) b(y, \eta) \Big|_{y=x, \eta=\xi}$$

- ▶ **(Formal) Adjoint.**

$$Op_h^W(a)^* = Op_h^W(\bar{a}).$$

Theorem (Calderón-Vaillancourt). There are $C > 0$ and $N \geq 0$ such that for all $a \in S(1, g_0)$ and all $h \in (0, 1]$

$$\|Op_h^W(a)\|_{L^2 \rightarrow L^2} \leq C \max_{|\alpha+\beta| \leq N} \|\partial_x^\alpha \partial_\xi^\beta a\|_{L^\infty}.$$

Ψ DO (part 2): The Gårding inequality

Theorem (Gårding inequality).

- ▶ **(Semiclassical version).** If $a \in S(1, g_0)$ satisfies $a \geq 0$, then there exists $C > 0$ such that

$$\operatorname{Re}(Op_h(a)u, u) \geq -Ch\|u\|^2, \quad h \in (0, 1], u \in S.$$

- ▶ **(“Classical” version)** If $a \in S^1 = S(\langle \xi \rangle, g_{\text{class}})$ satisfies $a \geq 0$, then there exists $C > 0$ such that

$$\operatorname{Re}(Op_1(a)u, u) \geq -C\|u\|^2, \quad u \in S.$$

Remark. Works for systems of Ψ DO too. In the purely scalar case, there is the Fefferman-Phong inequality ($h \rightarrow h^2$).

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Ψ DO (part 2): Proof of the Gårding inequality (semiclassical version)

For $(q, p) \in \mathbb{R}^{2n}$, define the (unitary) **phase space translation operator**

$$U(q, p)f(x) = e^{ix \cdot p} f(x - q).$$

Pick then $\eta \in \mathcal{S}$ such that $\|\eta\|_{L^2} = 1$ and define the **wave packets** $(\eta_{q,p})_{(q,p) \in \mathbb{R}^{2n}}$

$$\eta_{q,p}(x) = U(q, p)\eta(x).$$

Define the Bargmann transform

$$(Bu)(q, p) = (\eta_{q,p}, u)_{L^2}, \quad u \in \mathcal{S}.$$

Lemma. B maps $\mathcal{S}(\mathbb{R}^n)$ into $\mathcal{S}(\mathbb{R}^{2n})$.

One also has

$$(B^*\Psi)(x) = \iint \Psi(q, p)\eta_{q,p}(x)dqdp.$$

Proposition (resolution of identity). $B^*B = (2\pi)^n I.$

Ψ DO (part 2): Proof of the Gårding inequality (semiclassical version)

The resolution of identity $B^*B = (2\pi)^n I$ allows to define the **anti-Wick quantization**: given $a \in L^\infty(\mathbb{R}^{2n})$, we let M_a be the multiplication operator by a , and

$$Op^{aW}(a) = (2\pi)^{-n} B^* M_a B.$$

Using that

$$\|(2\pi)^{-n/2} B u\|_{L^2(\mathbb{R}^{2n})} = \|u\|_{L^2(\mathbb{R}^n)},$$

one has

$$\|Op^{aW}(a)\|_{L^2 \rightarrow L^2} \leq \|a\|_{L^\infty}.$$

The main property is the obvious fact that

$$a \geq 0 \quad \implies \quad Op^{aW}(a) \geq 0.$$

Important observation: all this holds for *any* profil η .

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Ψ DO (part 2): Proof of the Gårding inequality (semiclassical version)

Assume that η is *even* and *real valued* (e.g. Gaussian) and consider

$$\eta_h(x) = h^{-\frac{n}{4}} \eta(x/h^{\frac{1}{2}}), \quad \mathcal{W}_h(x, \xi) = (2\pi h)^{-n} e^{ix \cdot \xi} \eta(x/h^{\frac{1}{2}}) \hat{\eta}(\xi/h^{\frac{1}{2}})$$

(\mathcal{W}_h is normalized in $L^1(\mathbb{R}^{2n})$). Then

Proposition. $\mathcal{O}p^{aW}(a) = \mathcal{O}p_h(a * \mathcal{W}_h)$

Proof of the theorem.

$$\begin{aligned} \mathcal{O}p_h(a) &= \mathcal{O}p^{aW}(a) - \mathcal{O}p^{aW}(a) + \mathcal{O}p_h(a) \\ &= \underbrace{\mathcal{O}p^{aW}(a)}_{\geq 0} + \underbrace{\mathcal{O}p_h(a - a * \mathcal{W}_h)}_{=O_{L^2 \rightarrow L^2}(h)} \end{aligned}$$

using (for the $O(h)$) the Calderón-Vaillancourt theorem and the fact that

$$\|\partial_x^\alpha \partial_\xi^\beta (a - a * \mathcal{W}_h)\|_{L^\infty} \lesssim h.$$

Ψ DO (part 2): The Beals Theorem

We have already mentioned the (exact) formulas

$$[x_j, Op_h(a)] = ihOp_h(\partial_{\xi_j} a), \quad [h\partial_j, Op_h(a)] = hOp_h(\partial_{x_j} a).$$

They still hold for the Weyl calculus. Setting $ad_B A := [B, A]$ then, if $A_h = Op_h^W(a)$ for some $a \in S(1, g_0)$, for any N and

$$B_k \in \{Op_h^W(x_1), \dots, Op_h^W(x_n), Op_h^W(\xi_1), \dots, Op_h^W(\xi_n)\}$$

one has

$$ad_{B_1} \cdots ad_{B_N} A_h = h^N Op_h^W(\underbrace{a_N}_{\in S(1, g_0)}) = O_{L^2 \rightarrow L^2}(h^N).$$

Theorem. Let $A_h : \mathcal{S} \rightarrow \mathcal{S}'$ be a family of continuous operators. Then $A_h = Op_h^W(a(h))$ for some bounded family $(a(h))_{h \in (0, 1]}$ of $S(1, g_0)$ iff for all N and all B_k as above

$$ad_{B_1} \cdots ad_{B_N} A_h = O_{L^2 \rightarrow L^2}(h^N).$$

An example of application of the Beals theorem

Proposition. Assume that m is an order function and that $m \in S(m, g_0)$. Then, if h is small enough, $Op_h(a)$ has an inverse which is a ΨDO , i.e. there exists a bounded family $(\tilde{m}_h)_{h \ll 1}$ in $S(1/m, g_0)$ such that

$$Op_h^W(\tilde{m}_h)Op_h^W(m) = I = Op_h(1).$$

Proof. Observe that $1/m$ is an order function. Compute then

$$Op_h^W(1/m)Op_h^W(m) = I + hOp_h^W(b(h)) =: A_h,$$

with $(b(h))_{h \in (0,1]}$ bdd in $S(1, g_0)$. By Calderón-Vaillancourt, A_h is invertible for $h \ll 1$, and $\|A_h^{-1}\| \rightarrow 1$ as $h \rightarrow 0$. By iteration of

$$[B_k, A_h^{-1}] = -A_h^{-1}[B_k, A_h]A_h^{-1} = O_{L^2 \rightarrow L^2}(h)$$

the Beals Theorem gives $A_h^{-1} = Op_h^W(a(h))$ for some $(a(h))_{h \in (0,1]}$ bounded in $S(1, g_0)$. We then get result with $Op_h^W(\tilde{m}_h) := Op_h^W(a(h))Op_h^W(1/m)$. \square

Quantum mechanical interpretation and some dynamics

Classical mechanics: particles living in \mathbb{R}_x^n with motion described in the **phase space** $\mathbb{R}^{2n} = \mathbb{R}_x^n \times \mathbb{R}_\xi^n$, e.g.

$$m\ddot{x}(t) = -\nabla V(x(t))$$

reads equivalently

$$\begin{cases} \dot{x}(t) &= \frac{\xi(t)}{m}, \\ \dot{\xi}(t) &= (\nabla V)(x(t)) \end{cases} \quad \text{or} \quad \frac{d}{dt}(x(t), \xi(t)) = X_H(x(t), \xi(t))$$

where X_H is the **Hamiltonian vector field** associated with the (energy) **Hamiltonian** H

$$H(x, \xi) = \frac{1}{2m}|\xi|^2 + V(x), \quad X_H = (\nabla_\xi H, -\nabla_x H) = \left(\frac{\xi}{m}, -\nabla V\right)$$

Quantum mechanical interpretation and some dynamics

We denote the (Hamiltonian) flow of X_H by Φ_H^t , i.e.

$$\Phi_H^t(y, \eta) = (x(t), \xi(t)) \quad \text{with} \quad \begin{cases} \frac{d}{dt}(x(t), \xi(t)) = X_H(x(t), \xi(t)) \\ (x(0), \xi(0)) = (y, \eta) \end{cases}$$

(we assume it is complete, e.g. if $\nabla^2 V$ bounded). Other interesting types of H include

$$H(x, \xi) = \sum_{j,k=1}^n g^{jk}(x) \xi_j \xi_k$$

with $(g^{jk}(x))$ the coefficients of the co-metric tensor associated to a Riemannian metric \rightarrow **geodesic flow**.

Quantum mechanical interpretation and some dynamics

Functions $a : \mathbb{R}^{2n} \rightarrow \mathbb{R}$ are called **classical observables**.

For instance

- ▶ $a(x, \xi) = x_j$ (position in the j -th direction)
- ▶ $a(x, \xi) = \xi_k$ (k -th component of the momentum)
- ▶ $a(x, \xi) = H(x, \xi)$ (total energy)
- ▶ many others...

One may want to “observe” such quantities along the dynamics, i.e. consider $a \circ \Phi_H^t$.

Quantum mechanical interpretation and some dynamics

Quantum mechanics: particles living in \mathbb{R}_x^n are now described by **wave functions** $\psi \in L^2(\mathbb{R}^n, \mathbb{C})$ and the motion by the **Schrödinger equation**

$$\begin{cases} i\hbar\partial_t\psi(t) &= \hat{H}_\hbar\psi(t) \\ \psi(0) &= \psi_0 \end{cases} \quad (\hbar \sim \text{Planck's constant})$$

where \hat{H}_\hbar , the **quantum Hamiltonian**, is a selfadjoint operator on L^2 e.g.

$$\hat{H}_\hbar = -\frac{\hbar^2}{2m}\Delta_x + V(x).$$

This operator is obtained by **quantizing** the classical Hamiltonian $H(x, \xi) = \frac{|\xi|^2}{2m} + V(x)$, i.e.

$$\hat{H}_\hbar = \text{Op}_\hbar(H).$$

Quantum mechanical interpretation and some dynamics

Quantum observables are now operators on L^2 , which can be obtained by quantizing classical observables=symbols, i.e.

$$\hat{A}_{\hbar} = Op_{\hbar}(a).$$

Ψ DO provide a way to quantize classical observables. Note that the **composition theorem** of Ψ DO gives

$$\begin{aligned} [Op_{\hbar}(a), Op_{\hbar}(b)] &= Op_{\hbar}(a)Op_{\hbar}(b) - Op_{\hbar}(b)Op_{\hbar}(a) \\ &= Op_{\hbar}\left(ab + \frac{\hbar}{i}\nabla_{\xi}a \cdot \nabla_x b + O(\hbar^2)\right) - \\ &\quad Op_{\hbar}\left(ba + \frac{\hbar}{i}\nabla_{\xi}b \cdot \nabla_x a + O(\hbar^2)\right) \\ &= \frac{\hbar}{i}Op_{\hbar}(\{a, b\}) + O(\hbar^2) \end{aligned}$$

where $\{a, b\}$ is the **Poisson bracket** of a and b . This is an aspect of the **uncertainty principle** of quantum mechanics.

Quantum mechanical interpretation and some dynamics

The quantum flow is described by the **unitary** group $U_{\hbar}(t) = e^{-i\frac{t}{\hbar}\hat{H}_{\hbar}}$ and we can “observe” quantities like

$$(U_{\hbar}(t)\psi_0, \hat{A}_{\hbar}U_{\hbar}(t)\psi_0) = (\psi_0, \underbrace{U_{\hbar}(-t)\hat{A}_{\hbar}U_{\hbar}(t)}_{\hat{A}_{\hbar}(t)}\psi_0).$$

Theorem. For $a \in C_0^{\infty}(\mathbb{R}^{2n})$ and reasonable V , if $\hat{A}_{\hbar} = Op_{\hbar}(a)$ then for all (fixed) t ,

$$\hat{A}_{\hbar}(t) = Op_{\hbar}(a \circ \Phi_H^t) + O_{L^2 \rightarrow L^2}(\hbar)$$

Interpretation: In the semiclassical limit $\hbar \rightarrow 0$, the Schrödinger equation converges to the **Liouville equation**

$$\partial_t a(t) - \{H, a(t)\} = \partial_t a(t) - X_H \cdot \nabla_{x,\xi} a(t) = 0.$$

Correspondence between the classical Φ_H^t and quantum $U_{\hbar}(t)$ flows (vs Brownian motion for heat like semigroups)

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Semiclassical measures

Assume that (\hbar_k) is a (positive) sequence going to 0 and that we are given a sequence (ψ_k) of normalized functions in L^2 . Define the sequence of distributions $(\mu_k) \in \mathcal{D}'(\mathbb{R}^{2n})$ by

$$\langle \mu_k, a \rangle := (\psi_k, Op_{\hbar_k}(a)\psi_k).$$

These are distributions since $|\langle \mu_k, a \rangle| \leq \|Op_{\hbar_k}(a)\|_{L^2 \rightarrow L^2}$ is controlled by some seminorm of a in $S(1, g_0)$.

Assume then that, up to possible extraction, μ_k converges to some μ .

Proposition. The limit μ is a probability measure on \mathbb{R}^{2n} (called **semiclassical measure** associated to (ψ_k)).

Proof. To get a measure, it suffices to show $a \geq 0 \Rightarrow \langle \mu, a \rangle \geq 0$. This follows from the Gårding inequality since, if $a \geq 0$,

$$\langle \mu_k, a \rangle \geq -C_a \hbar_k \|\psi_k\|^2 \rightarrow 0, \quad k \rightarrow \infty.$$

It is a probability measure since $\langle \mu_k, 1 \rangle = \|\psi_k\|^2 = 1$. \square

Invariant semiclassical measures

Assume that (ψ_k) is a sequence of (normalized) *eigenfunctions* of $\hat{H}_{\hbar_k} = \text{Op}_{\hbar_k}(H)$ (say with $H(x, \xi) = |\xi|^2/2m + V(x)$).

Proposition. Any semiclassical measure μ associated to (ψ_k) is invariant by the classical Hamiltonian flow:

$$\langle \mu, a \circ \Phi_H^t \rangle = \langle \mu, a \rangle, \text{ for all } a \text{ and } t.$$

Proof. Up to extraction

$$\begin{aligned} \langle \mu, a \rangle &= \lim_{k \rightarrow \infty} (\psi_k, \text{Op}_{\hbar_k}(a)\psi_k) \\ &= \lim_{k \rightarrow \infty} (U_{\hbar_k}(t)\psi_k, \text{Op}_{\hbar_k}(a)U_{\hbar_k}(t)\psi_k) \\ &= \lim_{k \rightarrow \infty} (\psi_k, U_{\hbar_k}(-t)\text{Op}_{\hbar_k}(a)U_{\hbar_k}(t)\psi_k) \\ &= \lim_{k \rightarrow \infty} (\psi_k, \text{Op}_{\hbar_k}(a \circ \Phi_H^t)\psi_k) + O(\hbar_k) = \langle \mu, a \circ \Phi_H^t \rangle \quad \square \end{aligned}$$

Interest: describe how/where eigenfunctions are “localized” in the classical phase space.